

- [13] G. U. Lee, L. A. Chrisey, R. J. Colton, *Science* **1994**, 266, 771–773.
- [14] E.-L. Florin, V. T. Moy, H. E. Gaub, *Science* **1994**, 264, 415–417.
- [15] V. T. Moy, E.-L. Florin, H. E. Gaub, *Science* **1994**, 266, 257–259.
- [16] G. U. Lee, D. A. Kidwell, R. J. Colton, *Langmuir* **1994**, 10, 354–357.
- [17] Y. Harada, M. Kuroda, A. Ishida, *Langmuir* **2000**, 16, 708–715.
- [18] U. Dammer, O. Popescu, P. Wagner, D. Anselmetti, H.-J. Güntherodt, G. N. Misevic, *Science* **1995**, 267, 1173–1175.
- [19] W. Baumgartner, P. Hinterdorfer, W. Ness, A. Raab, D. Vestweber, H. Schindler, D. Drenckhahn, *Proc. Natl. Acad. Sci. USA* **2000**, 97, 4005–4010.
- [20] R. R. Schmidt, K.-H. Jung in *Preparative Carbohydrate Chemistry* (Ed.: S. Hanessian), Marcel Dekker, New York, **1997**.
- [21] The experiments were performed with a PicoSPM atomic force microscope, from Molecular Imaging, in a Teflon liquid cell. Commercial silicon nitride cantilevers (Olympus) were used for this study. They were rectangular shaped cantilever, 200  $\mu\text{m}$  long, with a 0.05  $\text{Nm}^{-1}$  nominal spring constant. This real spring constant was calibrated according to the dynamic behavior of the cantilever and to the equipartition theorem.<sup>[29, 30]</sup> The tips were coated with a five nanometer thick layer of chromium and then a 25 nm thick layer of gold.
- [22] C. D. Bain, E. B. Troughton, Y.-T. Tao, J. Evall, G. M. Whitesides, R. G. Nuzzo, *J. Am. Chem. Soc.* **1989**, 111, 321–335.
- [23] R. U. Lemieux, *Acc. Chem. Res.* **1996**, 29, 373–380.
- [24] J. M. de la Fuente, A. G. Barrientos, T. C. Rojas, J. Rojo, J. Cañada, A. Fernández, S. Penadés, *Angew. Chem.* **2001**, 113, 2318–2321; *Angew. Chem. Int. Ed.* **2001**, 40, 2258–2261.
- [25] T. Boland, B. D. Ratner, *Proc. Natl. Acad. Sci. USA* **1995**, 92, 5297–5301.
- [26] P. F. Luckham, K. Smith, *Faraday Discuss.* **1998**, 111, 307–320.
- [27] A. Chilkoti, T. Boland, B. D. Ratner, P. S. Stayton, *Biophys. J.* **1995**, 69, 2125–2130.
- [28] G. I. Bell, *Science* **1978**, 200, 618–627.
- [29] J. L. Hutter, J. Bechhoefer, *Rev. Sci. Instrum.* **1993**, 64, 1868–1873.
- [30] D. A. Walters, J. P. Cleveland, N. H. Thomson, P. K. Hansma, M. A. Wendman, G. Gurley, V. Elings, *Rev. Sci. Instrum.* **1996**, 67, 3583–3590.

## The $\text{CH}_3\text{N}$ Diradical: Experimental and Theoretical Determinations of the Ionization Energies\*\*

Wang Jing, Sun Zheng, Zhu Xinjiang, Yang Xiaojun, Ge Maofa, and Wang Dianxun\*

Methylnitrene, the  $\text{CH}_3\text{N}$  diradical, is one of the most interesting of all organic free radicals, and it has drawn a lot of attention in both experiments<sup>[1–5]</sup> and theory.<sup>[6–11]</sup> As the simplest alkyl nitrene, the general class of molecules with the form  $\text{RN}$ ,  $\text{CH}_3\text{N}$  is a dielectronic radical similar to carbene

$\text{CH}_2$ ; and is thought to be an important intermediate in many organic and inorganic reactions.<sup>[12–14]</sup> Berry suggests<sup>[1]</sup> that the alkyl nitrene should have electron states similar to  $\text{NH}$  which has been well characterized by spectroscopists. However, much less is known about alkyl nitrenes such as  $\text{CH}_3\text{N}$ ,  $\text{CH}_3\text{CH}_2\text{N}$ ,  $(\text{CH}_3)_2\text{CHN}$ , or  $(\text{CH}_3)_3\text{CN}$ . There are no reports on the ionization energy of the nitrenes until now, but these ionization energies, especially the first ionization energy, are important physical and chemical data.

When combined with ab initio molecular orbital calculations, HeI photoelectron spectroscopy (PES) gives a lot of information on the electronic and vibrational structures, the configuration, the bonding characteristics, as well as the relative photoionization cross sections of different molecular orbitals. Of course, the most important information obtained from the PES study is the ionization energy of the different molecular orbitals. Pople<sup>[15]</sup> and Ng<sup>[16]</sup> and their co-workers pointed out that the ab initio Gaussian2 (G2) calculation can correctly predict the first ionization energy.

We have shown how to generate a continuous flowing beam of short-lived species by the microwave discharge or pyrolysis of the parent species, allowing us to perform PES studies on reactive open-shell species such as  $\text{NO}_3$ ,<sup>[17]</sup>  $(\text{CH}_3)_2\text{N}$ ,<sup>[18]</sup>  $\text{CH}_3\text{O}$ ,<sup>[19]</sup>  $\text{CH}_3\text{S}$ ,<sup>[19]</sup>  $\text{CH}_3\text{SS}$ ,<sup>[20]</sup>  $\text{CH}_3\text{CH}_2\text{S}$ ,<sup>[21]</sup>  $(\text{CH}_3)_3\text{CO}$ <sup>[22]</sup> and further demonstrated that the G2 calculation can also predict the ionization energies of different ionic states. Herein, we present the experimental and theoretical determinations of the ionization energies of different ionic states for the  $\text{CH}_3\text{N}$  diradical; this is also the first report on ionization energy for the nitrenes.

Calculations on the  $\text{CH}_3\text{N}_3$  molecule indicate that the bond between the  $\text{CH}_3\text{N}$  and  $\text{N}_2$  units is the weakest present and therefore pyrolysis of the  $\text{CH}_3\text{N}_3$  offers a simple route to the  $\text{CH}_3\text{N}$  diradical through loss of the  $\text{N}_2$  unit.<sup>[23–26]</sup> However, in the pyrolysis study of  $\text{CH}_3\text{N}_3$  carried out by Bock and co-workers,<sup>[24–26]</sup> no PES signals of the  $\text{CH}_3\text{N}$  diradical could be detected and PES signals attributed to  $\text{CH}_2\text{NH}$  appeared when  $\text{CH}_3\text{N}_3$  was heated to 770 K.

Through careful analysis of the experimental process used by Bock et al.<sup>[24–26]</sup> it is clear that: 1) the distance from the point of pyrolyzation of the parent compound to the photo-ionized point for the new species generated is too far (maybe more than half meter); 2) the pyrolysis temperature (770 K) for maintaining the new species is also too high. Obviously, it is difficulty to record the PES signals of transient species which have a short life-time, such as the  $\text{CH}_3\text{N}$  diradical. In fact, mixed PE spectra showing  $\text{CH}_3\text{N}_3$  and  $\text{CH}_2\text{NH}$  (these spectra are the same as Figure 1 reported by Bock and Ralph in ref. [26]) were always obtained from the pyrolysis of the  $\text{CH}_3\text{N}_3$  at  $650(\pm 0.5)^\circ\text{C}$  with the heater installed outside our PE spectrometer.<sup>[27]</sup> When the PES of the species generated by pyrolysis of  $\text{CH}_3\text{N}_3$  are recorded in situ (see Experimental Section) a lot of the short-lived active species is detected at a lower temperature in our PE spectrometer.

Figure 1 gives the PE spectrum of the product obtained from the pyrolysis of  $\text{CH}_3\text{N}_3$  at  $145(\pm 0.5)^\circ\text{C}$ , the spectrum is recorded in situ under protection of the  $\text{NO}$  species. Note-worthy is that the mixed PE spectra showing both  $\text{CH}_3\text{N}_3$  and  $\text{CH}_2\text{NH}$  are also obtained at  $215(\pm 0.5)^\circ\text{C}$  if  $\text{NO}$  is not

[\*] Dr. W. Dianxun, W. Jing, S. Zheng, Z. Xinjiang, Y. Xiaojun, G. Maofa  
The Center for Molecular Sciences  
Institute of Chemistry  
The Chinese Academy of Sciences, Beijing 100080 (PR China)  
Fax: (+86) 10-62559373  
E-mail: wangdx@xx1.icas.ac.cn

[\*\*] This project was supported by the National Natural Science Foundation of China (Contract No. 29973051 and No. 20073052). Wang Jing, Sun Zheng, and Yang Xiaojun thanks the Chinese Academy of Sciences for a scholarship during the period of this work. Also we thank Qian Ximei, Li Chunhui, and Zhu Xinjiang for their contributions to this work during their time in our group.

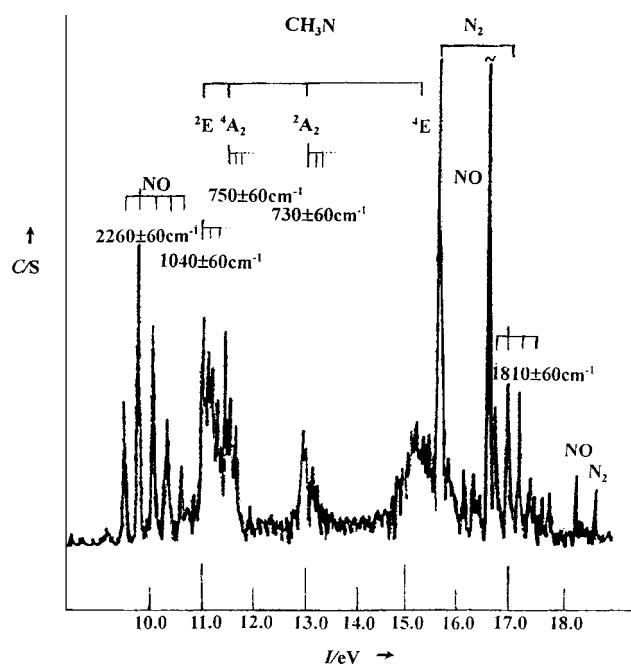


Figure 1. The PE spectrum (PES) of the  $\text{CH}_3\text{N}$  diradical obtained by pyrolysis of the  $\text{CH}_3\text{N}_3$  at  $145 (\pm 0.5)^\circ\text{C}$ .

present.<sup>[27]</sup> An expanded PE spectrum in which the fine vibrational structures could be clearly seen in the low ionization energy region ( $< 13.50$  eV) is given in Figure 2.

For calibration purposes it is worth noting that the PES bands of the NO molecule (bands at 9.54 eV, with fine vibrational structure ( $2260 \pm 60 \text{ cm}^{-1}$ ), at 16.56 eV (most intense band in Figure 1), and at 18.32 eV<sup>[28]</sup>), are clearly displayed in the PE spectra (Figures 1 and 2). The PES peaks at 15.60 eV, 16.98 eV with vibrational structure ( $1810 \pm 60 \text{ cm}^{-1}$ ), and 18.78 eV are characteristic of the  $\text{N}_2$  molecule.<sup>[28]</sup> Thus, the pyrolysis of the  $\text{CH}_3\text{N}_3$  generates  $\text{N}_2$ .

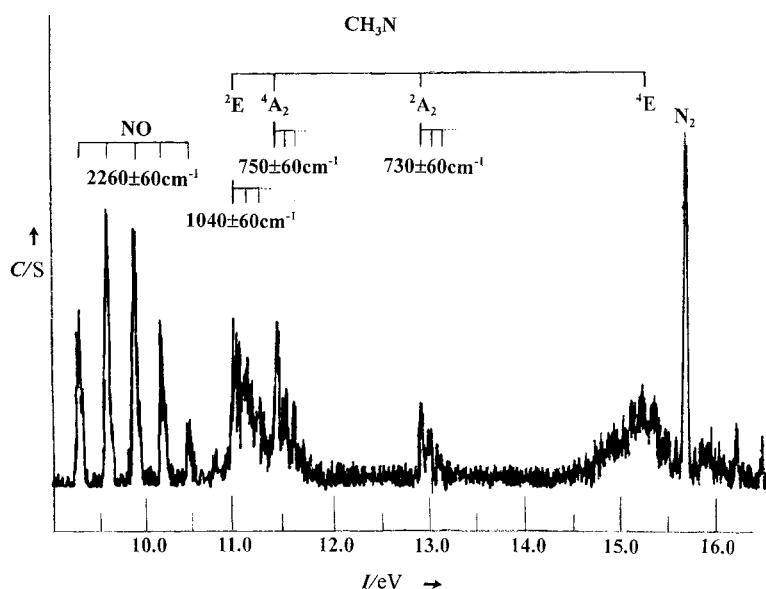


Figure 2. The expanded PE spectrum, in the lower ionization region ( $< 16.00$  eV), of the  $\text{CH}_3\text{N}$  diradical obtained by pyrolysis of the  $\text{CH}_3\text{N}_3$  at  $145 (\pm 0.5)^\circ\text{C}$ .

The PES bands at 11.04, 11.40, 12.85, and 15.53 eV are assigned to the  $\text{CH}_3\text{N}$  diradical. The peaks at 11.04, 11.40, and 12.85 eV are sharp with some structure to the higher ionization-energy side. This peak form is characteristic of the ionization of slightly bonding or slightly antibonding orbitals, i.e. the “vertical” ionization energy is equivalent to the adiabatic ionization energy of the PES band.<sup>[28]</sup> The peak at 11.04 eV is assigned to the lowest energy PES peak of the  $\text{CH}_3\text{N}$  radical cation, because the vibrational spacing ( $1040 \pm 60 \text{ cm}^{-1}$ ) is in excellent agreement with the reported vibrational frequency ( $1040 \text{ cm}^{-1}$ ) of the neutral  $\text{CH}_3\text{N}$  diradical,<sup>[3, 11]</sup> and the vibrational spacing of the PES band resulting from electronic ionization of a nonbonding orbital should be close to the vibrational frequency of the neutral species. The position of the first  $\text{CH}_3\text{N}$  PES peak matches well with the results of both G2 and DFT calculations, and is assigned as the  $^2\text{E}$  state of the  $\text{CH}_3\text{N}$  radical cation (see Table 1). According to the Frank–Condon principle, the

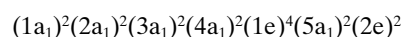
Table 1. The experimental vertical ionization energies ( $I_v$ ), calculated ionization energies ( $E_v$ ) according to  $\text{C}_{3v}$  symmetry from the G2 and DFT methods, and relative intensities of the PES signals observed from different ionic states of the  $\text{CH}_3\text{N}$  diradical.

PES experimental $I_v$ [eV] <sup>[a]</sup>	G2 energies $E_v$ [eV]	DFT energies $E_v$ [eV]	Cationic states	Relative intensity this work <sup>[b]</sup>	statistical ratio
11.04	11.063	11.006	$^2\text{E}^{[c]}$		
11.40	11.301	11.492	$^4\text{A}_2^{[d]}$	1.95	2
12.85	12.880	13.263	$^2\text{A}_2^{[d]}$	1.00	1
15.53	15.495	15.361	$^4\text{E}^{[e]}$		

[a]  $\pm 0.02$  eV. [b] Corrected for analyzer sensitivity  $\pm 0.02$ . [c] Ionic states come from electron ionization of the HOMO(2e) for  $\text{CH}_3\text{N}$  radical. [d] Ionic states come from electron ionization of the SHOMO( $5a_1$ ) for the  $\text{CH}_3\text{N}$  diradical. [e] Ionic states come from electron ionization of the 1e orbital of the  $\text{CH}_3\text{N}$  diradical.

energies of the ionic states are given using the vertical ionization energy ( $I_v$ ), chosen as the maximum of the PES bands, and the computed ionization energy ( $E_v$ ) should be obtained from the difference of the total energies of the resulting cation/neutral radical in  $\text{C}_{3v}$  symmetry.

The assignment of the PES bands for a new species can be supported by theoretical analysis. According to  $\text{C}_{3v}$  symmetry and a  $^3\text{A}_2$  ground state for the  $\text{CH}_3\text{N}$  diradical the molecular orbitals would be, in the order of increasing energy:



The removal of an electron from the HOMO will leave the ion in the singlet state  $^2\text{E}$ , whereas removal of an electron from each of the other valence shell orbitals in turn can leave both quartet and doublet states of the cation.

The HOMO(2e) of the  $\text{CH}_3\text{N}$  diradical has a dominant contribution from the lone-pair electrons of the nitrogen atom with a small degree of the C–N antibonding character. This means that removal of

an electron from the HOMO(2e) of CH<sub>3</sub>N, leading to the cation in the singlet state <sup>2</sup>E, should give a sharp PES peak with fine structure of 1040 ± 60 cm<sup>-1</sup> which is in excellent agreement with the reported vibrational frequency 1040 cm<sup>-1</sup> of the C–N stretch mode of the neutral CH<sub>3</sub>N diradical.<sup>[3, 11]</sup>

The orbital character of the 5a<sub>1</sub> of the CH<sub>3</sub>N diradical is essentially that of a lone-pair orbital of the nitrogen atom. This means that the removal of the electrons of the 5a<sub>1</sub> orbital should lead to two sharp peaks corresponding to <sup>4</sup>A<sub>2</sub> and <sup>2</sup>A<sub>2</sub> ionic states of the CH<sub>3</sub>N radical cation, and the intensity of the PES band for the <sup>4</sup>A<sub>2</sub> ionic state should be two times that of the <sup>2</sup>A<sub>2</sub> ionic state, because the intensity of the PES band represents the relative ionic statistical weights.

A careful analysis of the intensities of the two peaks at 11.40 and 12.85 eV for the PE spectrum of the pyrolyzed CH<sub>3</sub>N<sub>3</sub> shows that the intensity of the peak at 11.40 eV is about 1.95 times that of the peak at 12.85 eV. That is to say, the intensity ratio of the two peaks is in good agreement with the ratio of relative ionic statistical weight led by electron ionization of the 5a<sub>1</sub> orbital for the CH<sub>3</sub>N diradical, that is, they correspond to the <sup>4</sup>A<sub>2</sub> and <sup>2</sup>A<sub>2</sub> ionic states of the CH<sub>3</sub>N radical cation. The vibrational spacing (750 ± 60 cm<sup>-1</sup>) of the peak at 11.40 eV is very close to the vibrational spacing (730 ± 60 cm<sup>-1</sup>) of the peak at 12.85 eV. This fact also shows that the two PES peaks should arise from the result of electron ionization of the same orbital. The PES band centered near 15.53 eV could be attributed to removal of the electrons from the 1e orbital or maybe also includes the contribution of more orbitals in the deep valence shell of CH<sub>3</sub>N, because the band is very broad. As indicated in Table 1, PES experimental and both G2 and DFT theoretical results agree reasonably well.

### Experimental Section

Methyl azide CH<sub>3</sub>N<sub>3</sub> was prepared according to the methods reported in refs. [26, 29]. The purity of CH<sub>3</sub>N<sub>3</sub> was checked by mass spectrometry and NMR spectroscopy. The PE spectrum of the CH<sub>3</sub>N<sub>3</sub> obtained is the same as that reported<sup>[26]</sup> and did not vary from day to day before each experiment.

The PES experiment was performed in a double-chamber UPS Machine-II which was built specially, as described elsewhere,<sup>[17]</sup> to detect transient species. A continuous flowing CH<sub>3</sub>N diradical beam is produced in situ by pyrolysis of CH<sub>3</sub>N<sub>3</sub> at 145(±0.5) °C in a quartz tube by using a double heater inlet system, the heater was capable of producing 1200 °C. Thus, the new species was produced within 1–2 cm of the photoionization region of the spectrometer, it then passed through the photoionization point and was quickly pumped out. The double heaters were installed inside of the sample chamber. The PE spectrum of the CH<sub>3</sub>N diradical was measured at a resolution of about 30 meV as indicated by the Ar<sup>+</sup>(<sup>3</sup>P<sub>3/2</sub>) photoelectron band. Experimental vertical ionization energies (*E<sub>v</sub>* in eV) were calibrated by simultaneous addition of a small amount of NO to the sample, the first PES band with NO fine vibrational structure can be used for calibration the of the fine vibrational structure of the new species, NO can also stabilize the new species.<sup>[22]</sup>

Gaussian 2 and improved DFT calculations: to assign the PES bands of the CH<sub>3</sub>N diradical, both G2 and DFT calculations were performed on the <sup>3</sup>A<sub>2</sub> ground state of the CH<sub>3</sub>N diradical in the C<sub>3v</sub> symmetry and the ground and several ionic states of the CH<sub>3</sub>N radical cation. The G2 theoretical procedure has been described in detail by Curtiss et al.<sup>[15]</sup> The improved density functional theory (DFT) calculations are described in refs. [17–22]. The computed ionization energies (*E<sub>v</sub>* in eV) were obtained from the difference between the total energy of the resulting radical cation <sup>3</sup>A<sub>2</sub> ground state in the C<sub>3v</sub> symmetry to the total energy of the CH<sub>3</sub>N diradical.

Received: March 20, 2001 [Z16816]

- [1] R. S. Berry in *Nitrenes* (Ed.: W. Lwowski), Wiley-Interscience, New York, **1970**, p. 13.
- [2] C. Zetsch, F. Stuhl, *Chem. Phys. Lett.* **1975**, 33, 375.
- [3] E. L. Chappell, P. C. Engelking, *J. Chem. Phys.* **1988**, 89, 6007.
- [4] M. J. Travers, D. C. Cowles, E. P. Clifford, G. B. Ellison, *J. Am. Chem. Soc.* **1992**, 114, 8699.
- [5] D. A. Hrovat, E. E. Waali, W. T. Borden, *J. Am. Chem. Soc.* **1992**, 114, 8701.
- [6] M. S. Platz, *Acc. Chem. Res.* **1995**, 28, 487.
- [7] A. Gilles, J. Masanet, C. Vermeil, *Chem. Phys. Lett.* **1974**, 25, 345.
- [8] S. J. Kim, T. P. Hamilton, H. F. Schaefer III, *J. Am. Chem. Soc.* **1992**, 114, 5349.
- [9] J. H. Glowina, J. Misewich, P. P. Sorokin in *Supercontinuum Laser Sources* (Ed.: R. R. Alfano), Springer, New York, **1989**, p. 337.
- [10] C. Richards, C. Meredith, S. Kim, G. E. Quelch, H. F. Schaefer, *J. Chem. Phys.* **1994**, 100, 481.
- [11] M. J. Travers, D. C. Cowles, E. P. Clifford, G. B. Ellison, P. C. Engelking, *J. Chem. Phys.* **1999**, 111, 5349.
- [12] C. Wentrup, *Reactive Molecules: The Neutral Reactive Intermediates in Organic Chemistry*, Wiley, New York, **1984**, chap. 4.
- [13] R. G. Sadygov, D. R. Yarkony, *J. Chem. Phys.* **1997**, 107, 4994.
- [14] Y. Kurosaki, T. Takayanagi, K. Sato, S. Tsunashima, *J. Phys. Chem. A* **1998**, 102, 254.
- [15] L. A. Curtiss, K. Raghavachari, G. W. Trucks, J. A. Pople, *J. Chem. Phys.* **1991**, 94, 7221.
- [16] Z. X. Ma, C. L. Liao, C. Y. Ng, Y. S. Cheung, W. K. Li, T. Bear, *J. Chem. Phys.* **1994**, 100, 4870.
- [17] D. X. Wang, P. Jiang, X. M. Qian, G. Y. Hong, *J. Chem. Phys.* **1997**, 106, 22.
- [18] C. H. Qiao, G. Y. Hong, D. X. Wang, *J. Phys. Chem. A* **1999**, 103, 1972.
- [19] X. J. Zhu, M. F. Ge, J. Wang, Z. Sun, D. X. Wang, *Angew. Chem.* **2000**, 112, 2016; *Angew. Chem. Int. Ed.* **2000**, 39, 1940.
- [20] M. F. Ge, J. Wang, Z. Sun, D. X. Wang, *J. Chem. Phys.* **2001**, 114, 3051.
- [21] M. F. Ge, J. Wang, X. J. Zhu, Z. Sun, D. X. Wang, *J. Chem. Phys.* **2000**, 113, 1866.
- [22] Z. Sun, S. X. Zheng, J. Wang, D. X. Wang, *Chem. Eur. J.*, in press.
- [23] J. A. Leemakers, *J. Am. Chem. Soc.* **1933**, 55, 3098.
- [24] H. Bock, R. Dammel, *Angew. Chem.* **1987**, 99, 518; *Angew. Chem. Int. Ed. Engl.* **1987**, 26, 504.
- [25] H. Bock, *Polyhedron* **1988**, 7, 2429.
- [26] H. Bock, R. Dammel, *J. Am. Chem. Soc.* **1988**, 110, 5261.
- [27] X. M. Qian, Master thesis, **1997**; C. H. Li, Master thesis, **1998**; X. J. Zhu, Master thesis, **2000**, Institute of Chemistry, Chinese Academy of Science, Beijing, China.
- [28] K. Kimura, S. Katsumata, Y. Achiba, T. Yamazaki, S. Iwata, *Handbook of HeI Photoelectron Spectra of Fundamental Organic Molecules*, Japan Scientific Societies, Tokyo, **1980**.
- [29] O. R. Francis, J. G. Chester, *J. Phys. Chem.* **1957**, 61, 830.